

**₁ Radar reflectivity as a proxy for convective mass
₂ transport**

G. L. Mullendore,¹ A. J. Homann,¹ K. Bevers,² C. Schumacher,²

G. L. Mullendore, Department of Atmospheric Sciences, University of North Dakota, 4149
University Avenue, Stop 9006, Grand Forks, ND 58202, USA. (gretchen@atmos.und.edu)

¹Department of Atmospheric Sciences,
University of North Dakota, Grand Forks,
North Dakota

²Department of Atmospheric Sciences,
Texas A&M University, College Station,
Texas

3 **Abstract.** More observations of vertical mass transport in deep convec-
4 tion are needed to improve dynamical understanding of detrainment processes
5 and for verification of transport models. A methodology for using radar re-
6 flectivity as a direct observation of vertical transport of mass from the bound-
7 ary layer to the upper troposphere and lower stratosphere is investigated and
8 the “level of maximum detrainment” (LMD) is proposed. The case investi-
9 gated is the 26 January 1999 squall line from the TRMM-LBA field cam-
10 paign. Echo-top heights and dual-Doppler derived divergence profiles are used
11 to define the mass detrainment range. Over 10% of anvil echo tops occurred
12 above the sounding-derived LNB of 15.4 km during the mature stage of the
13 storm, and convective tops reached above 18 km. Anvil ice water content,
14 with a simple correction for ice fall-speed, is found to be a good proxy for
15 both the LMD, which for the storm analyzed is 11.25 km, and for the de-
16 trainment range of 6 to 17 km. More cases need to be analyzed to confirm
17 the strength of this methodology, but the case study presented shows a strong
18 correlation between anvil properties determined from radar reflectivity and
19 the mass detrainment profile. Thus, radar reflectivity can be used as an in-
20 dicator of the LMD to test model convective and transport parameteriza-
21 tions.

1. Introduction

22 Rapid vertical transport of chemical tracers from the boundary layer to the upper
23 troposphere and lower stratosphere (UTLS) is accomplished primarily through convective
24 transport in storms (e.g. Dickerson [1987]). In order for atmospheric chemical models to
25 properly simulate the UTLS, the models first must properly simulate the level at which
26 mass is detraining from deep convective updrafts. Better understanding of deep convective
27 transport is important not only for chemical models, but also dynamical models, as updraft
28 size and extent relates closely to the problem of heating profiles and water vapor transport,
29 crucial for global momentum budgets and radiative budgets, as highlighted in recent
30 studies (e.g. Schumacher et al. [2004]; Alexander et al. [2004]). Deep convective mass
31 transport has been simulated (e.g. Stenchikov et al. [1996]; Mullendore et al. [2005];
32 Barth et al. [2007]; Charboureau et al. [2007]), but without more observations, results
33 from cloud-resolving transport models remain relatively unconstrained.

34 The simplest way to estimate the level at which detrainment will occur is to calculate
35 the level of neutral buoyancy (LNB) from a sounding using parcel theory. The radiosonde
36 used, however, is often located a substantial distance in time or space from the immediate
37 pre-storm environment, and therefore may have sampled an air mass very different from
38 the air parcels being convected. Also, this parcel method of determining detrainment
39 level is highly idealized, not accounting for the variability in initial parcel conditions or
40 entrainment of environmental air along the upward trajectory of the surface parcel. The
41 entrainment amounts and profiles can be approximated (e.g. Emanuel [1991]; Kain and
42 Fritsch [1990]), but entrainment itself is highly variable, depending on factors such as the

43 local environmental profile, storm classification and storm size (Cohen [2000]; Mullendore
44 et al. [2005]).

45 Case studies have used chemical tracer retrievals from aircraft and sounding measure-
46 ments to analyze observations of tracer transport in vigorous deep convection (e.g. Poulida
47 et al. [1996]; Strom et al. [1999]; Hegglin et al. [2004]), but these observations are limited
48 both spatially and temporally. Satellite observations have provided data over large areas
49 on tracer plumes originating from deep convection (e.g. Ricaud et al. [2007]; Jiang et al.
50 [2007]), but in a particular satellite observational data set, either the horizontal or vertical
51 resolution of the measurements is coarse relative to the cloud-resolving regional models.
52 Indirect methods, such as inferring deep convective transport based on the radiative bal-
53 ance in the upper troposphere [Folkins and Martin, 2005], also lack the resolution for
54 cloud-resolving models.

55 To constrain cloud-resolving transport models, it is desirable to have a direct, cloud-
56 scale measurement of the heights at which storm outflow *actually occurred*. Mullendore
57 et al. [2005] found good agreement between the radar and model-simulated reflectivity
58 in the forward anvil structure associated with a Severe Thunderstorm Electrification and
59 Precipitation Study (STEPS) supercell (see their figure 14). The authors asserted that
60 the reflectivity in the anvil region shows the transport of hydrometeors and serves as a
61 proxy for tracer mass transport in deep convective updrafts. In this study, the goal is
62 to expand on this assertion using reflectivity and dual-Doppler derived velocities from a
63 squall line observed during the Tropical Rainfall Measuring Mission Large-Scale Biosphere-
64 Atmosphere (TRMM-LBA) campaign.

65 The TRMM-LBA field experiment was conducted in Brazil in January and February
66 of 1999 with the goals of learning more about tropical continental precipitation and vali-
67 dating TRMM satellite products [Silva Dias et al., 2002]. Cifelli et al. [2002] investigated
68 the differences in various storm properties occurring in different meteorological regimes of
69 TRMM-LBA, including storm morphology, rainfall amounts, and mass fluxes. This study
70 expands upon the use of dual-Doppler derived velocities to estimate upward transport
71 amounts integrated over the storm and proposes a methodology that could be used in
72 the absence of dual-Doppler analysis. The term “level of maximum detrainment” (LMD)
73 will be used to describe the level where maximum mass detrainment occurs. Cloud-top
74 retrievals provide an upper bound on LMD and indicate the altitudes where troposphere-
75 to-stratosphere exchange is most likely occurring, while reflectivity volumes illuminate the
76 details of the vertical mass detrainment profile.

77 The current study builds on the dual-Doppler analysis done by Cifelli et al. [2002] by
78 focusing on the deep convective mass transport, comparing the sounding-derived LNB to
79 the LMD, and investigating the possibility of using anvil reflectivity retrievals as a proxy
80 for convective mass transport. Section 2 gives a brief overview of the observations used.
81 Section 3 describes the analyses performed and presents the echo-top and mass transport
82 statistics. Summary and conclusions are presented in Section 4.

2. Observations

83 The observational platforms for TRMM-LBA included four radiosonde sites and two
84 radars, the NASA TOGA C-band radar and the NCAR S-pol radar (figure 1). The dual-
85 Doppler velocities are derived from the radial velocities observed by the NASA TOGA
86 radar and the NCAR S-pol radar. The radar quality control and dual-Doppler processing

87 was performed by researchers at Colorado State University [Cifelli et al., 2002; Lang and
88 Rutledge, 2002]. Horizontal velocities (U and V) were corrected for precipitation fall
89 speed by using a reflectivity-terminal fall velocity ($Z - V_t$) relationship. Vertical velocity
90 was derived from the horizontal divergence calculations with downward integration using
91 an exponential density profile. Note that although the horizontal divergence is available
92 from the original dual-Doppler analysis, this study uses the derived vertical velocity data
93 in order to look at upward-only velocities as well as calculate statistics on the vertical
94 velocities.

95 The specific case being studied is the storm that moved westward through the TRMM-
96 LBA observational domain on January 26th, 1999. The squall line formed at an outflow
97 boundary from previous convection several hundred kilometers northeast of the TRMM-
98 LBA sampling domain [Cifelli et al., 2002]. This storm was present in the region covered
99 by the dual-Doppler lobes between 1950 and 2210 UTC and had a classic leading line-
100 trailing stratiform structure [Houze, 1993]. The leading convective line and forward anvil
101 were present in the dual-Doppler region between 1950 and 2030 UTC, which was the time
102 period during which the storm was intensifying. The remaining times sampled primarily
103 the stratiform region of the squall line.

3. Results

104 Due to its location in relation to the approaching squall line the 18 UTC sounding
105 at Rebio-Jaru (figure 2, black line) was chosen as most representative of the pre-storm
106 environment, and was used to estimate the LNB; both graphical and numerical methods of
107 lifting a surface parcel returned a LNB altitude of 15.4 km. In contrast, the LNB calculated
108 using the Abracos Hill sounding (figure 2, gray line) was 13.8 km. This high variability

109 over a relatively small spatial area is due to the high variability in the boundary layer
110 readings, highlighting the difficulty in relying on sounding data for storm characteristics.

3.1. Echo Top Retrievals

111 Following the methodology described in Steiner et al. [1995], the raining regions of the
112 storm were classified as either convective or stratiform. The anvil was defined as the
113 reflectivity observations with echo tops at or above 6 km with no reflectivity at 2.5 km
114 following Frederick and Schumacher [2008].¹ The separation of the storm into convective
115 and stratiform rain regions has a long tradition in the literature, but the anvil regions
116 (i.e., horizontally homogeneous echo observed aloft but not reaching the ground) are often
117 ignored or lumped in with the stratiform rain classification when analyzing reflectivity
118 data. The anvil is, in fact, dynamically and radiatively distinct and is important for mass
119 transport considerations. Note that often all cirrus is categorized as anvil; here the anvil is
120 specifically the thick cirrus that is still attached to or recently produced by its convective
121 source. For the specific squall line analyzed in this study, anvil region is the forward anvil
122 and stratiform region is the trailing stratiform.

123 The storm was partitioned into three time periods; intensifying (1950-2030 UTC), ma-
124 ture (2040-2110 UTC), and dissipating (2120-2210 UTC). The echo tops, defined as the 0
125 dBZ line, for each region of the storm were then placed into altitude bins; the frequency
126 of echo tops occurring at a given altitude and time is shown in figure 3. Echo tops were
127 consistently produced in the region of the LNB (15.4 km) and although some of the ob-
128 served tops may represent negatively buoyant air, some transport is expected to occur at
129 all heights that hydrometeors penetrate, via turbulent mixing. In the intensifying stage
130 of the storm, hydrometeors in the stratiform rain region are reaching above the LNB and

131 strong convective updrafts at mid- to upper-levels are carrying hydrometeors to 18 km,
132 indicating the possibility of mixing at high altitudes and transport directly into the trop-
133 ical stratosphere. Echo tops for each storm region are highest in the mature stage of the
134 storm, indicating stronger and more persistent updrafts. In the dissipating stage of the
135 storm, vertical velocities are weaker and most hydrometeors are found below 16 km.

136 Mullendore et al. [2005] demonstrated that while turbulent processes do allow some
137 parcels to achieve neutral buoyancy at high altitudes by mixing with stratospheric air,
138 many parcels at cloud top in deep convective turrets are, in fact, negatively buoyant and
139 will subsequently descend. In contrast, the anvil region, on convective time scales of an
140 hour or two, is moving more horizontally than vertically and indicates a neutrally buoyant
141 air mass. Interestingly, the echo tops of the anvil region remain high even in the later
142 stage of the storm (figure 3c), showing that much of the anvil is generated by the deep
143 cells during the mature stage of the storm and then remains aloft.

3.2. Mass Transport Profiles

144 While echo-top heights provide an upper bound of deep convective mass transport, echo-
145 top data is insufficient to obtain a vertical profile and learn at what altitude the maximum
146 detrainment is occurring. The final extent of the anvil associated with a particular storm
147 depends on several factors not related to the strength and duration of the convective
148 updrafts, including radiative interactions, upper level wind shear and magnitude, and
149 upper level relative humidity. However, in an active storm, when time scales range from
150 30 minutes to several hours and spatial scales are on the order of the updraft width,
151 the dominant factors in anvil production will be divergence in the updraft (strength of
152 updraft) and upper level winds. While strong upper level winds can create some anvil

153 from overshooting hydrometeors (Garrett et al. [2004]; Wang [2003]), the majority of anvil
154 near the updraft will be due to vertical mass convergence. The extent and magnitude of
155 the convergence, and likewise, the extent and concentration of hydrometeors in the anvil,
156 describe the LMD.

157 **3.2.1. Single Cross-Section**

158 The convective line is redefined in this section as the region 25 km in width that is
159 bounded on the west by the first reflectivity points that extend from 2 km to at least 8 km
160 in altitude over a minimum of 5 km distance in the horizontal (figure 4a). This convective
161 region definition differs from the traditional Steiner et al. [1995] reflectivity definition and
162 from the methodology used in the previous section, and was chosen so that all regions of
163 vertical mass convergence that may affect the forward anvil formation, including weaker
164 dissipating cells, would be included. Note that although the stratiform region of the
165 storm is responsible for some amount of mass transport, much of the transport occurring
166 in the stratiform region is transport of recycled air that originated in the convective
167 region [Gamache and Houze, 1983], or relatively small magnitude vertical transport of
168 environmental air from the mid-troposphere [Gamache and Houze, 1982]; the likelihood
169 of venting of boundary layer air with subsequent transport to high altitudes is unlikely
170 in the stratiform region. Hence, this study focuses on the deep vertical transport in the
171 convective regions of the storm. The possible role stratiform transport plays in the further
172 lifting of air already processed by the deep convective updrafts will be investigated in a
173 separate study. The convectively generated anvil is defined as in the previous section but
174 constrained to be within 20 km of the convective line.

175 Figure 4c shows the mean vertical velocity (gray line) and mean reflectivity (black line)
 176 in the convective region of the cross-section shown. Values above 16 km are less represen-
 177 tative of a mean value, as the number of echoes sharply decreases at that altitude. The
 178 velocity peaks at approximately 11 km, suggesting that the majority of parcels experience
 179 negative buoyancy above that height and hence begin to slow down. This slice was chosen
 180 for its well-developed anvil. The contour plot of the vertical velocity (panel b) suggests
 181 that the updraft sampled is dissipating, with a maturing cell that has moved farther back
 182 from the leading edge and is now elevated and cut-off from surface air.

Figure 4d shows the total divergence ($\frac{d\rho w}{dz}$, gray line) in the convective region and the
 total ice water mass (black line) in the anvil region of the cross-section. To calculate the
 vertical mass divergence, ρw is calculated at each velocity point (resolution is 0.5 km in
 the vertical and 1.0 km in the horizontal) by approximating $\rho(z)$ as $(\rho_0 \exp^{z/H})$, where
 $\rho_0 = 1.22 \text{ kg/m}^3$ and $H = 7 \text{ km}$. Then a first order difference was used to calculate $\frac{d\rho w}{dz}$
 along each column. To calculate the total vertical mass divergence profile, the divergence is
 integrated horizontally at each vertical level. The ice mass was estimated from reflectivity
 by assuming the majority of hydrometeors in the forward anvil are ice, then using equation
 3 from Leary and Houze [1979]:

$$I = 8.0 \times 10^{-3} Z_I^{0.61}$$

183 where I is the ice water content (IWC), and Z_I is the reflectivity obtained after converting
 184 from the effective reflectivity by adding 6.7 dB. Note that although a range of magnitudes
 185 for IWC could be generated by choosing a different relation, the shape of the IWC profile
 186 would not change significantly.

187 The peak in negative vertical mass divergence indicates a peak in horizontal detrainment
188 near 12 km, due to decreasing velocity in the updraft located about 50 km from SPOL
189 (panel b). This primary peak in detrainment corresponds well with the peak in the
190 total anvil ice water mass. A secondary peak in detrainment near 14 km corresponds to
191 the elevated maximum located near 50 km from SPOL. As described above, the updraft
192 sample is not at its most vigorous stage; the elevated maximum seen in the vertical
193 velocity contours is partially responsible for the notch in divergence at 13 km. The notch
194 is also partially caused by positive divergence of downdrafts in the volume. The notch
195 is decreased in magnitude somewhat (i.e., convergence increases; figure 4d, dashed line)
196 by calculating velocities from positive vertical velocities only (figure 4c, dashed line). As
197 the goal is to capture detrainment primarily caused by the deep convective updrafts, only
198 positive vertical velocities will be used to calculate deep convective transport in subsequent
199 plots.

200 The correlation between vertical updraft mass convergence and anvil ice water mass
201 supports the theory that anvil mass is a good proxy for convective detrainment level.
202 However, not all cross-sections show such a strong relationship, as the updraft may be
203 dissipated, or the updraft is strong but so young that the anvil has not yet had time
204 to mature. To truly understand the overall correlation, a more statistical picture of the
205 storm structure is needed.

206 **3.2.2. Transport Statistics**

207 To investigate the more general relationship between mass detrainment and anvil
208 heights, contoured frequency by altitude diagrams (CFADs, Yuter and Houze [1995])
209 of the anvil and convective regions are calculated for multiple time steps (figure 5, panels

210 a-d). The CFADs for reflectivity, vertical velocity, and vertical divergence ($\frac{d\rho w}{dz}$) in the
211 convective line (figure 5, panels a-c) are constructed using 4 dBZ-sized bins centered from
212 2 dBZ to 58 dBZ, 2 m s⁻¹-sized bins centered from -20.5 m s⁻¹ to 20.5 m s⁻¹, and 2x10⁻³
213 kg m⁻³ s⁻¹-sized bins centered from -0.02 to 0.02 kg m⁻³ s⁻¹, respectively. The CFAD
214 for ice water content in the forward anvil (figure 5d) is constructed using 0.01 g m⁻³-sized
215 bins centered from 0.01 g m⁻³ to 0.2 g m⁻³. The outermost bins extend to capture any
216 values outside the range specified. The contours shown are 0.1, 1, 10, and 30% frequency.
217 This analysis only includes times 1950 UTC to 2030 UTC, as afterwards, the leading
218 convective line was no longer clearly defined.

219 Although we have used a different method to define the convective region, the CFADs
220 of convective reflectivity and convective velocity, as well as the mean reflectivity (figure
221 5e) and mean vertical velocity (figure 5f), compare closely with the CFAD results shown
222 in Cifelli et al. [2002] (see their figures 8 and 9). To determine the detrainment levels
223 and the LMD, the vertical divergence in the convective region was calculated. The tilt
224 in the 30% frequency contours (figure 5c) shows more vertical divergence, and horizontal
225 entrainment, occurring at the low altitudes and more vertical convergence, and horizontal
226 detrainment, occurring at the high altitudes. The profile of total divergence at each
227 altitude (figure 5g) shows that there is net entrainment below 5 km and net detrainment
228 between 5 km and 16.5 km. The altitude of maximum detrainment, as determined from
229 maximum vertical convergence of positive vertical velocities, is 11.25 km. Panel h shows
230 the anvil ice mass at each time snapshot, with the total anvil ice mass at time 2010 UTC
231 in bold. Time 2010 UTC has the maximum ice mass as the forward anvil region within
232 20 km of the convective line starts leaving the domain at time 2020 UTC. The total ice

233 water mass peaks at 10 km, with over 75% of the detrainment occurring at 2010 UTC
234 between 7.5 and 13.0 km and over 50% between 8.5 and 11.5 km. The maximum in the ice
235 water content at 10 km matches fairly well with the detrainment profile shown in panel
236 g. However, ice is expected to fall as the anvil advects away from the source region of the
237 updraft, necessitating a possible correction to the anvil altitudes.

238 **3.2.3. Ice Fall Speed Adjustment**

239 Assuming an ice fall speed of 0.3 m s^{-1} [Houze, 1993] and a mean outflow speed, we
240 can introduce a correction to the anvil heights. The mean outflow speed is the speed the
241 anvil is being transported westward in relation to the convective line. Over the analysis
242 times, the average westward speed of the line is calculated to be 11 m s^{-1} , and the average
243 dual-Doppler east-west velocity in the anvil is 16 m s^{-1} westward. Using the mean outflow
244 speed of 5 m s^{-1} , the anvil would displace downwards 60 m for every 1 km in the horizontal
245 it advects away from the convective line. Because the anvil has been constrained to the
246 area within 20 km of the convective line, the maximum vertical adjustment is $1.2 \text{ km} +/-$
247 0.4 km (assuming 0.1 m s^{-1} uncertainty in the ice fall speed).

248 Figure 6a shows the new IWC CFAD and ice detrainment profile for time 2010 UTC
249 using the fall speed adjustment. This adjustment has lifted the anvil estimated maximum
250 detrainment level to 10.75 km (figure 6b), which compares well with the detrainment
251 profile maximum of 11.25 km estimated from the total vertical divergence (shown again
252 in figure 6c). For the case studied here, adjusted anvil location is a good proxy for mass
253 detrainment profile.

254 The final LMD of approximately 11.25 km is 4 km below the Rebio-Jaru LNB of 15.4
255 km and 2.5 km below the LNB of 13.8 km from the Abracos Hill sounding. Simulations

256 of this squall line are needed to investigate the mechanisms at work in more detail. The
257 detrainment profile obtained from the analyses presented here will provide constraints on
258 those future simulations.

4. Summary and Conclusions

259 The January 26th squall line from the TRMM-LBA field campaign is used as a test
260 case for a methodology developed to use radar reflectivity as a direct observation of
261 vertical transport of mass from the boundary layer to the upper troposphere and lower
262 stratosphere. The “level of maximum detrainment” (LMD) is proposed to clarify the
263 difference between the actual detrainment profile in deep convection and the level of
264 neutral buoyancy (LNB) obtained from parcel theory.

265 Echo-top heights occur in a range containing the LNB, demonstrating the upper bound
266 on convective transport and the altitudes at which turbulent mixing may cause irreversible
267 transport. Over 10% of anvil echo tops occur above the LNB during the mature stage of
268 the storm, and convective echo tops reach above 18 km. As some transport is expected at
269 all heights that hydrometeors penetrate, this suggests some air is convectively transported
270 directly into the stratosphere. By the dissipating stage of the storm, almost 10% of the
271 anvil echo tops are still above the LNB. An algorithm for separating out the regions of
272 cloud extent from Geostationary Operational Environmental Satellites (GOES) visible
273 and infrared images that represent young, convectively-produced anvil based on TRMM
274 radar observations is currently being investigated.

275 The reflectivity in the storm’s forward anvil region is tested and found to be a good
276 proxy for the LMD. The anvil region of interest is the anvil formed directly by convective
277 turrets, in this storm defined as reflectivity retrievals with echo tops at or above 6 km with

278 no reflectivity at 2.5 km that occur west (forward) and within 20 km of the convective line.
279 A vertical profile of total anvil ice water mass is obtained by converting reflectivity to ice
280 water mixing ratio and applying a simple ice fall-speed correction depending on horizontal
281 distance from the convection. The resulting ice water mass profile correlates well with
282 the LMD determined from the vertical velocity divergence profile. In the storm analyzed,
283 the divergence-calculated LMD is at 11.25 km and estimated at 10.75 km using anvil
284 reflectivity alone, with a detrainment range from 6 to 17 km, while the LNB, calculated
285 with parcel theory, is calculated to be 15.4 km. More cases need to be analyzed to
286 confirm the strength of this methodology, but this case study shows a strong correlation
287 between anvil properties and the mass detrainment profile. Research is currently being
288 conducted to determine how to best apply this methodology to different storm types (e.g.
289 non-squall multicell storms, supercells). Also, tests are being conducted to use NEXRAD
290 observations. Although some additional error will be introduced when storm-relative anvil
291 speeds are not available, this methodology still shows promise for identifying the LMD.
292 Radar reflectivity can be used as an indicator of the LMD to test model convective and
293 transport parameterizations and also to constrain models that resolve convection.

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Notes

1. The gridded reflectivity field contained poorly resolved features in the most northern 20 km of the domain and most eastern 5 km, so these regions were excluded from the analyses.

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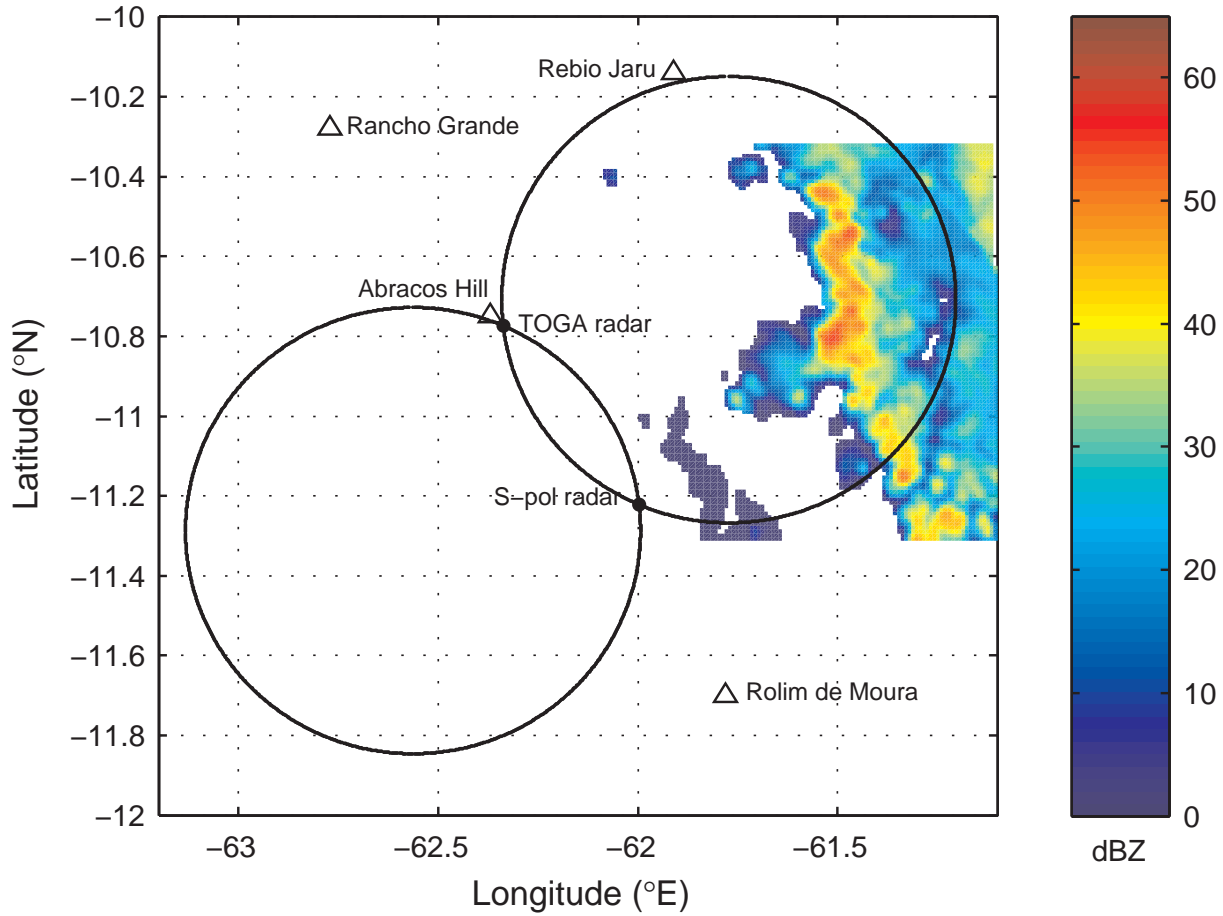


Figure 1. The TRMM-LBA sounding (triangles) and radar platforms (filled circles) used in this study. The dual-Doppler lobes are shown as open circles. Reflectivity is shown for altitude of 5 km at time 1950 UTC.

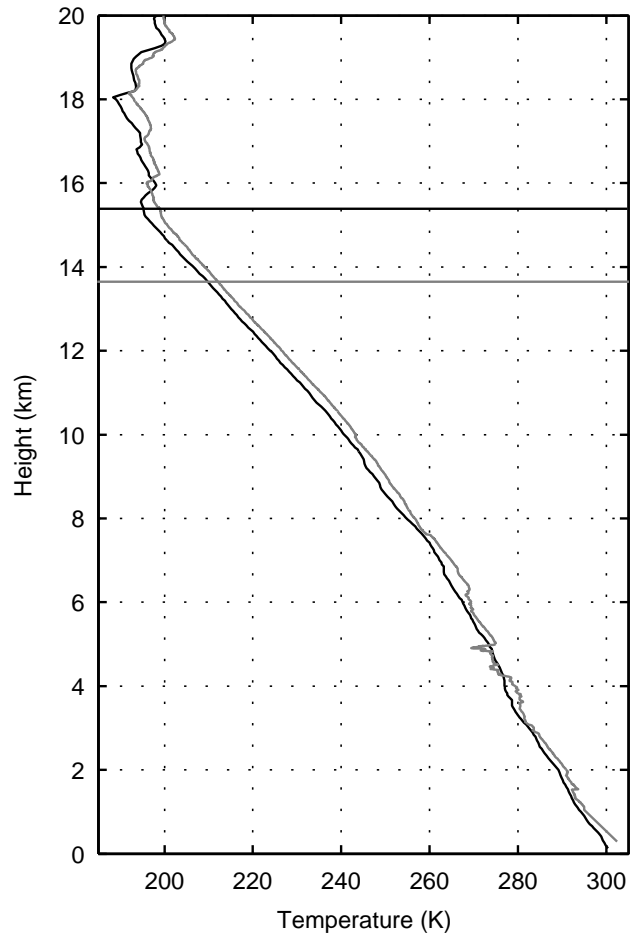


Figure 2. Temperature profiles and LNB altitudes (horizontal lines) for the January 26, 1999, 18 UTC soundings from Rebio-Jaru (black) and Abracos Hill (gray).

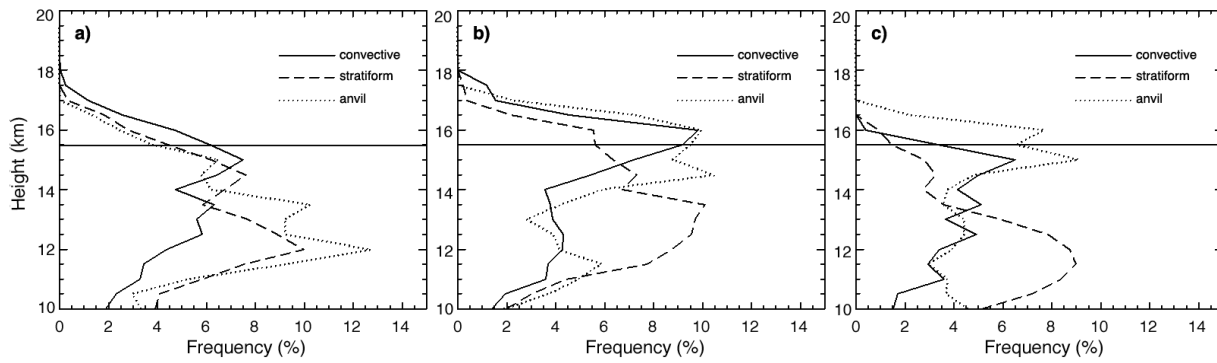


Figure 3. Reflectivity echo tops from the a) intensifying, b) mature, and c) dissipating stages of the January 26th squall line. Frequency of occurrence of echo top at each altitude is shown for convective (solid line), stratiform (dashed line), and anvil (dotted line) regions of the storm. The LNB is shown as a horizontal solid line.

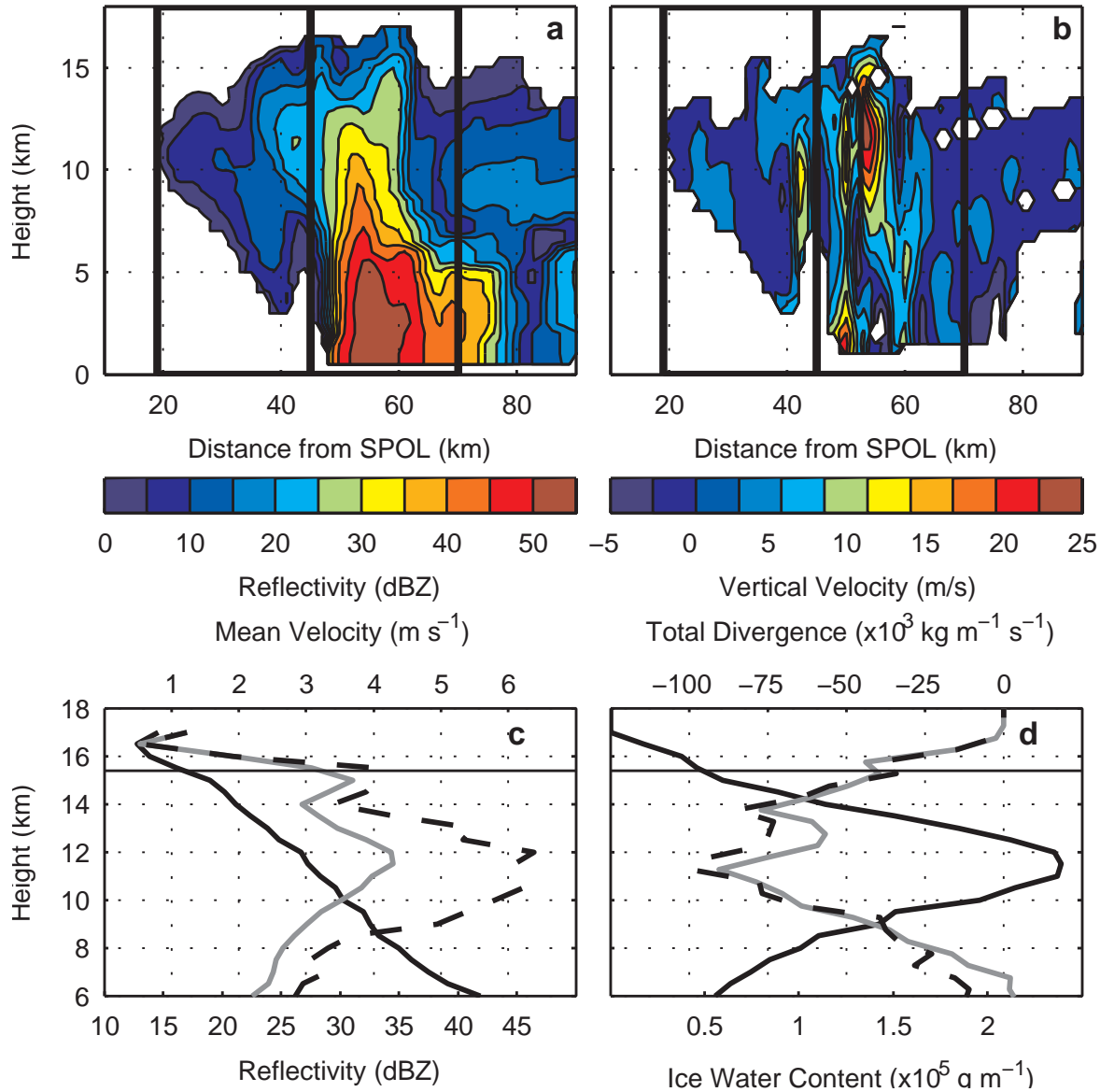


Figure 4. (a, b) Vertical cross-section taken perpendicular to line at 1950 UTC, showing reflectivity (a) and vertical velocity (b) in x and z, with the anvil and convective regions shown as boxes. The dual-Doppler domain extends an additional 60 km to the west. (c) Mean vertical velocity (gray line), mean vertical positive velocity (dashed line) and mean reflectivity (black line) in the convective region of the cross-section. (d) Total divergence ($\frac{d\bar{w}}{dz}$, gray line) and total divergence of positive velocities (dashed line) in the convective region and the total ice water mass (black line) in the anvil region of the cross-section. Horizontal line shows Rebio-Jaru LNB.

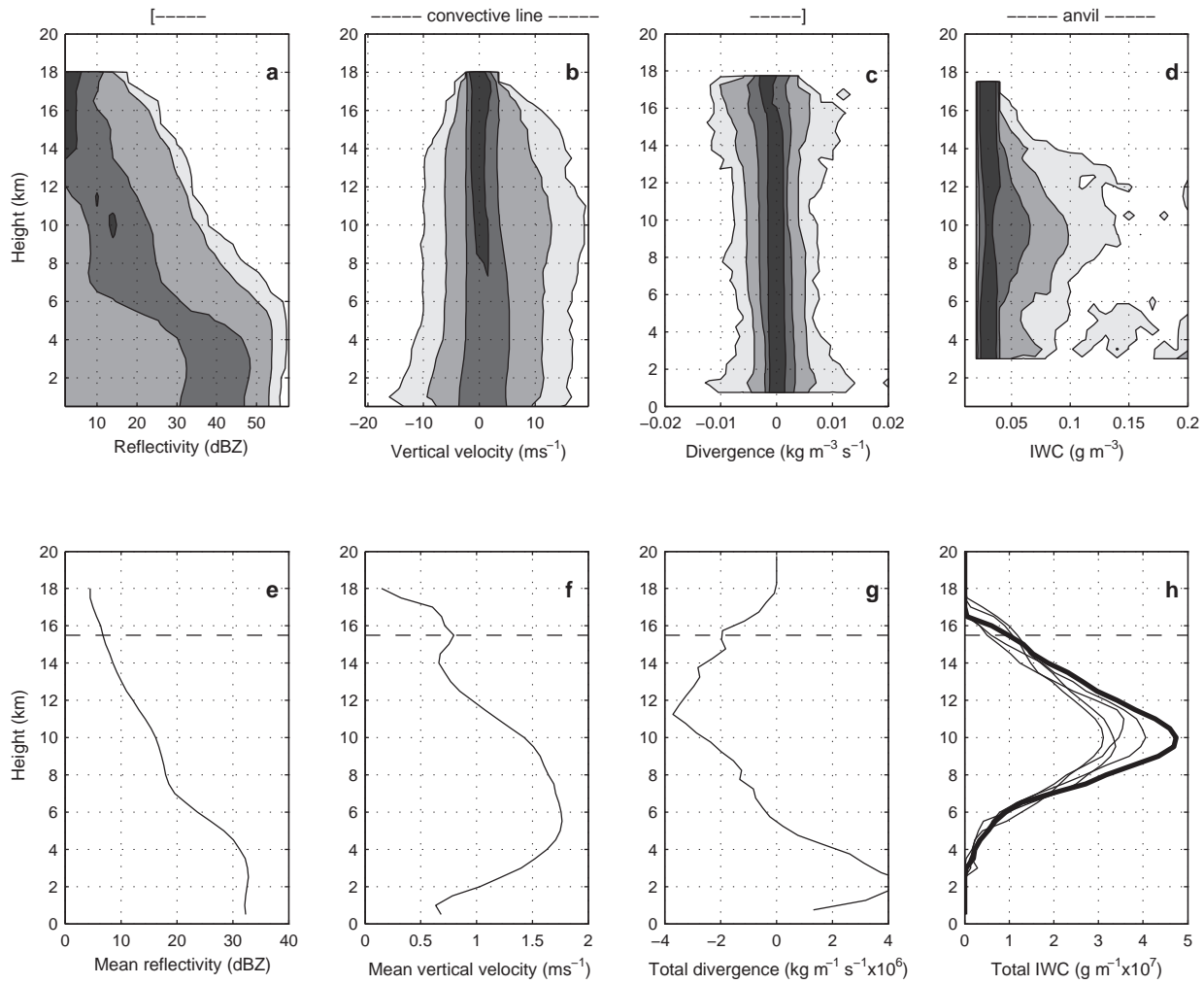


Figure 5. Analysis of times 1950 UTC to 2030 UTC in the Jan. 26 squall line. Panels a-c: Contoured frequency by altitude diagrams (CFADs) of a) reflectivity, b) vertical velocity, and c) vertical divergence in the convective line. Panel d: CFAD of ice water content in the forward anvil. Contours shown are 0.1, 1, 10, and 30% frequency. See text for details on bins chosen. Panels e-f: Mean e) reflectivity and f) vertical velocity for the convective line region at each altitude. Panel g: Total vertical divergence in convective line at each altitude. Panel h: Total ice water mass in forward anvil at each altitude, at each time slice. The anvil at 2010 UTC is shown as a thick solid line. In all bottom panels, the LNB is shown as a horizontal dashed line.

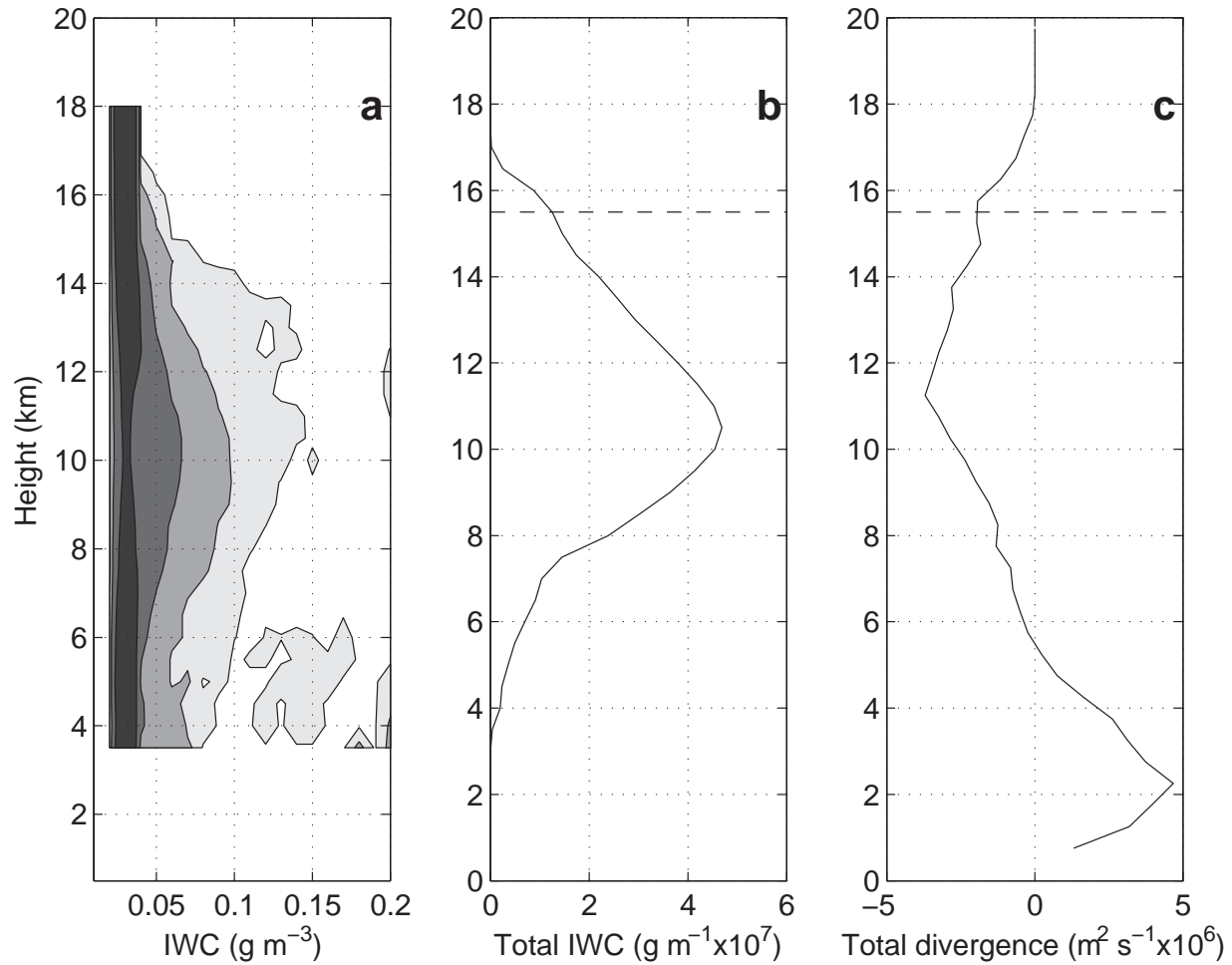


Figure 6. a) CFAD of ice water content from 1950-2030 UTC and b) total ice water mass at 2010 UTC for anvil profiles adjusted by ice fall speed. c) Total vertical divergence for positive vertical velocities. LNB shown as dashed horizontal line in panels b and c.